

Six-Inch High-Purity Lead Halide Perovskite Wafer Derived from Ceramic Manufacturing Technique

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Lead halide perovskites exhibit exceptional optoelectronic properties but face industrialization barriers due to the inability to fabricate large-area, high-quality wafers. Inspired by ceramic manufacturing techniques, a 6-inch high-purity perovskite wafer is developed, achieving carrier mobility, lifetime, and defect concentrations comparable to single crystals. This method demonstrates universality across diverse perovskites and enables heterojunction wafers, marking significant progress in carrier dynamics control. As a result, an X-ray sensing array with 256×256 pixels is constructed using a $10 \times 10 \text{ cm}^2$ perovskite heterojunction wafer, which exhibits a sensitivity of $36532 \mu\text{CGy}_{\text{air}}^{-1} \text{cm}^{-2}$ and a low detection limit of $139 \text{ nGy}_{\text{air}} \text{ s}^{-1}$, superior to those in a single-crystal detector ($10640 \mu\text{CGy}_{\text{air}}^{-1} \text{cm}^{-2}$ and $247 \text{ nGy}_{\text{air}} \text{ s}^{-1}$). This breakthrough establishes a scalable pathway to industrial-scale perovskite optoelectronics, overcoming critical manufacturing barriers while enabling high-performance radiation imaging systems through wafer-level heterostructure engineering.

1. Introduction

The unique combination of low-cost raw materials, facile synthesis, and tailored optical properties has propelled lead halide perovskites (APbX_3 ; $\text{A}^+ = \text{CH}_3\text{NH}_3^+$ [MA^+], $\text{CH}(\text{NH}_2)_2^+$ [FA^+], Cs^+ ; $\text{X}^- = \text{I}^-$, Br^- , Cl^-) into diverse applications spanning solar cells,^[1] direct X-ray/ γ -ray detectors,^[2] scintillators,^[3] light-emitting diodes,^[4] lasers,^[5] nonlinear optics,^[6] and visible/infrared photodetectors.^[7] Despite these promising attributes, perovskite-based device development remains at the lab-scale stage, and industrial manufacturing still faces numerous difficulties.

Taking silicon-based devices as an example, they rely on standardized wafer-scale Si substrates (6–12-inch diameters) compatible with automated micro/nanofabrication

workflows. However, the manufacture of perovskite wafers is still unsatisfactory. For example, melt-based techniques like Bridgman growth can produce all-inorganic variants (e.g., CsPbX_3),^[8] but organic-inorganic hybrid perovskites suffer from thermal decomposition of volatile organic cations.^[9] Similarly, solution-based crystallization methods, though versatile, exhibit poor reproducibility and scalability. Currently, state-of-the-art perovskite single crystals rarely exceed 5 cm in lateral dimensions and frequently contain structural inhomogeneities that compromise purity and optical quality.^[9,10] These limitations bring a critical mismatch with the silicon-based device preparation technique. The absence of standardized perovskite wafer manufacturing hinders the compatibility with existing semiconductor foundry infrastructure—a prerequisite for transitioning from laboratory prototypes to commercial-scale production.

Recent advances in polycrystalline ceramic processing offer a promising pathway to fabricate large-area perovskite wafers. Lead halide perovskite wafers, synthesized via hot-pressing or sintering, have demonstrated X-ray detection sensitivities surpassing those of single-crystal MAPbI_3 ^[11] and rivaling commercial CdTe detectors.^[12] For instance, $\text{Cs}_3\text{Bi}_2\text{I}_3$ wafers^[13] exhibit high carrier mobility and radiation hardness, while $\text{Cs}_2\text{AgBiBr}_6$ ^[14] double perovskites achieve low dark currents and enhanced stability. However, existing ceramic fabrication methods yield opaque, mechanically fragile materials with limited thickness and poor

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optical transparency—properties directly linked to grain boundary defects, secondary phases, and porosity.^[12a,15]

Taking transparency as an example, it requires dense, phase-pure microstructures with minimized light scattering—a hallmark of high-quality wafers. Transparent perovskite wafers can facilitate multifunctional applications including sensing arrays, nonlinear optics, rare-earth-doped lasers, and monolithic heterojunctions. Recent studies on lead-free variants like (BA)₂CsAgBiBr₇^[16] and (CH₃NH₃)₃BiI₉^[15d-f] highlight the growing interest in combining apparent transparency and optoelectronic performance. Nevertheless, no work has yet demonstrated freestanding, large-scale perovskite wafers (> 5 cm in diameter) with sub-millimeter thickness and visible/infrared transparency.

In this article, we demonstrate a high-purity, 6-inch, pinhole-free transparent perovskite wafer with ceramic manufacturing techniques. The high-quality perovskite wafers exhibit excellent optical and electrical properties, including defect levels, carrier mobility, and carrier lifetime comparable to those of single crystals. The multi-factor interaction mechanism of a precise control strategy within ceramic manufacturing techniques for improving the quality of perovskite wafers is revealed. The fabrication method can be expanded to different A-site cations and different halogen perovskites, as well as multilayer perovskite heterojunctions. Finally, an X-ray sensing array with 256 × 256 pixels was established on our perovskite heterojunction wafer with a high sensitivity of 36532 μCGyair⁻¹ cm⁻² and a low detection limit of 139 nGyair s⁻¹ superior to MAPbBr₃ single-crystal detector (10640 μCGyair⁻¹ cm⁻² and 247 nGyair s⁻¹). These results confirm the potential of our perovskite transparent wafer for constructing functional commercial devices using industrial manufacturing techniques.

2. Results and Discussion

In general, the high optical transparency of perovskite single crystals indicates their excellent properties, including phase-purity, crystallinity, defect, and optoelectronic behavior.^[17] Opacity in polycrystalline materials mainly originates from optical anisotropy of grains, impurities or secondary phases at grain boundaries, high surface roughness ($\gg \lambda/10$), and pores.^[18] Achieving high optical transparency in thick polycrystalline wafers requires stringent control of microstructural properties, including porosity, and crystal grain stacking. The schematic diagram in **Figure 1a** shows the difference in perovskite grain between the normal perovskite tablet with the hot-pressing technique and our proposed transparent perovskite wafer (MAPbBr₃). It is reasonable to believe that perovskite grain growth or grain boundary fusion is important to produce the pinhole-free perovskite wafer and ensure good transparency.

To achieve the transparent perovskite wafer, we employ a ceramic hot-pressing technique with a precise control strategy. In brief, we selected small perovskite single crystals from the solution process and fine-grind them to a micrometer scale. Subsequently, a perovskite wafer with a thickness of 0.6 mm can be fabricated via a ceramic fabrication method. **Figure S1** (Supporting Information) illustrates this process. **Figure S2** (Supporting Information) displays the influence of temperature, pressure, and time during hot-pressing on MAPbBr₃ wafers' transparency. Interestingly, MAPbBr₃ wafers change from opaque to transparent

by adjusting the temperature, pressure, and time during the fabrication. With the accurate temperature of 150 °C and pressure of 200 MPa for 60 min, the obtained MAPbBr₃ wafer is transparent, which looks like a MAPbBr₃ single crystal. **Figure 1b** depicts the photograph of the ShunGui Laboratory Building (Southeast University) captured through a MAPbBr₃ transparent wafer, demonstrating the exceptional photon transmittance properties.

Surprisingly, the as-fabricated transparent MAPbBr₃ wafers can reach 6-inch wafer-grade. **Table S1** (Supporting Information) summarizes the reported perovskite single crystals and polycrystalline tablets. To our best knowledge, our transparent MAPbBr₃ wafer is the largest and its size is significantly larger than that of perovskite single crystals. **Figure 1c,d** shows its optical image in the lab. Notably, one side is patterned with a gold electrode array. The gold electrode array is clearly visible on the reverse side. **Figure 1e** presents the optical transmission spectrum. Our MAPbBr₃ wafer exhibits a high optical transparency of $\approx 68\%$ comparable to that of the single crystal (75%) within the 700–1100 nm range,^[19] which is significantly higher than that in normal MAPbBr₃ tablets.

To verify that the transparency leads to excellent optoelectronic properties in the MAPbBr₃ wafer, we employed the time-of-flight technique to acquire the charge carrier drift time (τ_r) under different applied biases^[20] (**Figure S3**, Supporting Information). The hole and electron mobility in MAPbBr₃ transparent wafer were calculated from the transit time and wafer thickness, yielding values of 99 and 180 cm² V⁻¹ s⁻¹, respectively (**Figure 1f,g**). The value is comparable to that of MAPbBr₃ single crystals (107 and 198 cm² V⁻¹ s⁻¹), which is significantly higher than that of normal MAPbBr₃ tablet (76 and 119 cm² V⁻¹ s⁻¹). This result indicates that our MAPbBr₃ transparent wafer manufacturing technique can efficiently eliminate the grain boundary and facilitate grain growth or fusion.

Quantification of wafer densities using the Archimedes method, as depicted in **Figure 2a** and **Table S2** (Supporting Information), confirms a progressive increase in density with pressure up to 200 MPa, with the density approaching that of a MAPbBr₃ single crystal. However, when the pressure further goes up to 300 MPa, the density shows a clear decrease, which can be ascribed to the internal voids from wafer fragmentation. It should be noted that the time and temperature remain constant (150 °C and 60 min) throughout the process. If the time is shorter, incomplete compaction and fusion of MAPbBr₃ grains result in opacity at the wafer edges. Conversely, prolonging the time to 120 min leads to powder adherence to the mold, causing surface fracturing during demolding.

To verify the grain growth in our ceramic manufacture, we conducted a high-resolution X-ray rocking curve analysis of the (100) diffraction peak (**Figure 2b**; **Figure S4**, Supporting Information). According to the Scherrer formula, the grain size is inversely proportional to the full width at half maximum (FWHM). The results confirm that the wafer prepared at 150 °C exhibits larger grain sizes and better crystalline quality compared to those prepared at 50, 70, and 180 °C. Higher temperatures, such as 180 °C, induce the vaporization of the organic component MABr, leading to grain breakdown. To further accurately design the perovskite wafer manufacturing, Statistical Package for the Social Sciences software was used to generate the orthogonal table of temperature, pressure, and time. As shown in **Tables S3** and **S4**

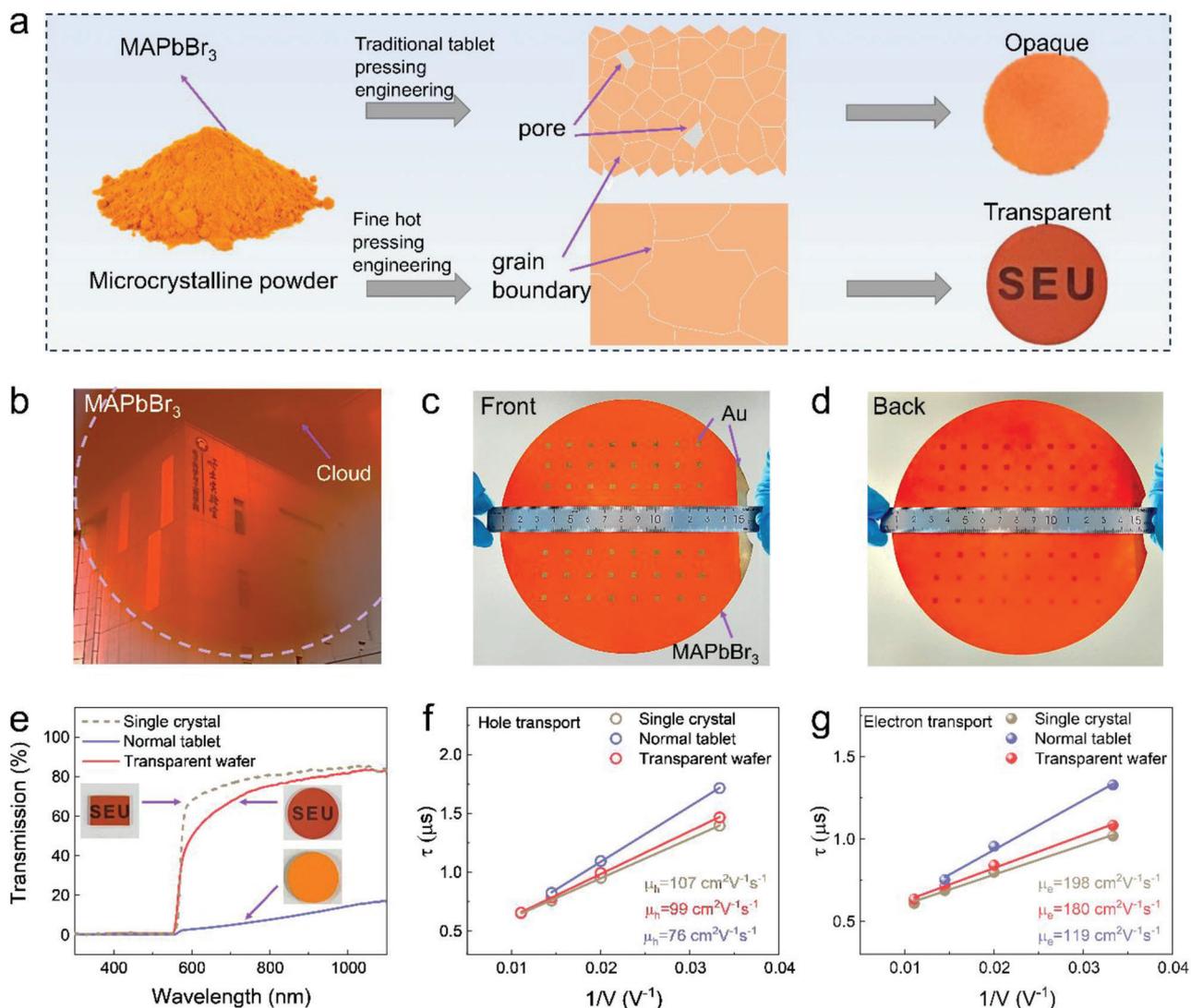


Figure 1. Six-inch transparent MAPbBr₃ wafers. a) The schematic illustration of the differences in grain boundary between the transparent wafer and the normal tablet. b) An image of the ShunGui Laboratory Building (Southeast University) captured through a MAPbBr₃ transparent wafer. c) Photograph of the 6-inch transparent wafer with gold electrode array on its front surface. d) Photograph of the 6-inch transparent wafer with the gold electrode array on the backside. e) The optical transmission spectrum. f) Hole and g) electron mobility.

(Supporting Information), the minimum K arithmetic mean values indicate the optimal settings are K₃ (150 °C), K₂ (200 MPa), and K₃ (60 min). Scanning electron microscopy images, as depicted in Figure S5 (Supporting Information), reveal that stress reduces the gaps between the grains, further confirming grain boundary fusion. Figure S6 (Supporting Information) depicts the dark *I*-*V* curves of Au/MAPbBr₃/Au devices. For each specific factor, the transparent MAPbBr₃ wafer prepared at 150 °C, 200 MPa, and 60 min exhibits the lowest dark current.

Figure 2c presents surface and cross-sectional Scanning electron microscope (SEM) images of single crystals and wafers with a thickness of $\approx 600 \mu\text{m}$. The transparent wafer exhibits clear fused grain boundaries, demonstrating a continuous and compact surface morphology. In contrast, the normal tablet surface contains obvious pores. Cross-sectional SEM analysis further proves the compact density in the transparent wafer, as well as

the single crystal. Figure 2d presents time-resolved photoluminescence (TRPL) spectra. The results demonstrate that the carrier lifetime of transparent MAPbBr₃ wafer significantly increases to 2.6 μs , which is quite longer than that of a normal tablet and approaches the value of 3.17 μs in a single crystal. As shown in Figure 2e, the $\mu\tau$ of MAPbBr₃ transparent wafer ($1.8 \times 10^{-4} \text{ cm}^2 \text{V}^{-1}$) is much higher than that of the normal tablet ($4.3 \times 10^{-5} \text{ cm}^2 \text{V}^{-1}$), comparable to that of single crystal ($2.3 \times 10^{-4} \text{ cm}^2 \text{V}^{-1}$).

Halide vacancies at grain boundaries generally act as ion migration channels in lead-based perovskites.^[21] Temperature-dependent conductivity measurements were conducted to evaluate the electronic and ionic conductivity in transparent MAPbBr₃ wafers. The activation energy (E_a) of ion migration can be determined using the Nernst–Einstein relation:

$$\sigma(T) = \left(\frac{\sigma_0}{T} \right) \exp\left(-\frac{E_a}{Tk_B} \right) \quad (1)$$

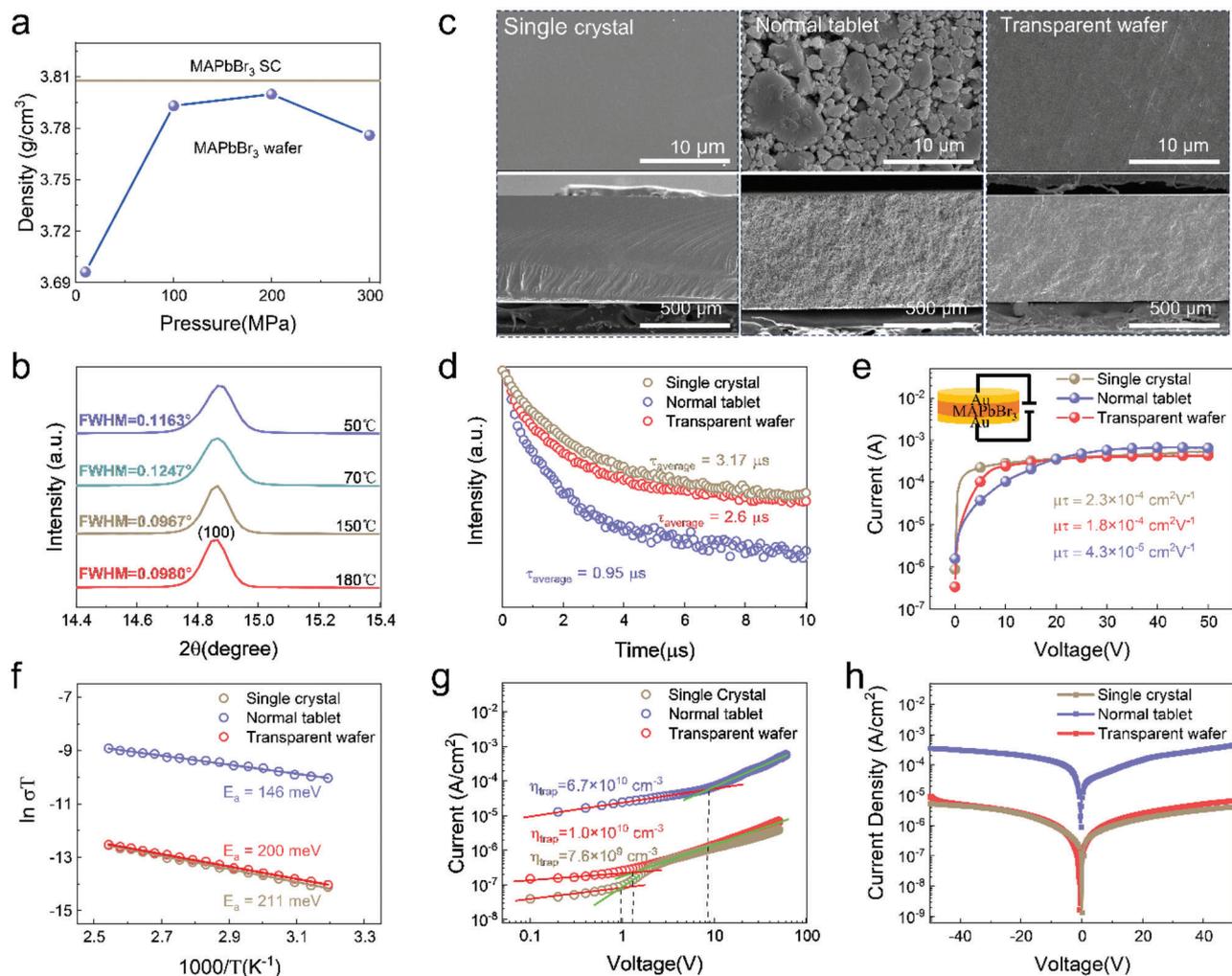


Figure 2. Characterization of the MAPbBr₃ transparent wafer. a) Density of MAPbBr₃ wafers prepared under different pressures. b) Magnified view of the (200) diffraction peak. c) SEM images of MAPbBr₃ single crystal and wafers. d) Time-resolved PL spectra. e) Bias-dependent photoconductivity. f) Temperature-dependent conductivity. g) SCLC curves of single crystals and wafers. h) *I*–*V* curve of Au/MAPbBr₃/Au devices based on single crystal and wafers under dark.

where σ_0 is a constant, T is the temperature, and k_B is the Boltzmann constant. As shown in Figure 2f, the E_a in transparent wafer (200 meV) is much higher than that of a normal tablet (146 meV), approaching the level of a single crystal (211 meV). It indicates the suppressed ion migration. As well, the dark current baseline was examined under a 30 V voltage for ≈ 7 min (Figure S7, Supporting Information) to extract the dark current drift D :^[22]

$$D = \left| (J_t - J_0) / tE \right| \quad (2)$$

where J_0 and J_t are the dark current densities at the beginning and end points, respectively, t is the duration time, and E is the electric field. The current baseline of a transparent wafer is much more stable with D of 5.6×10^{-4} nA cm⁻¹ s⁻¹ V⁻¹, which is three orders of magnitude lower than that of the normal tablet (1.73×10^{-1} nA cm⁻¹ s⁻¹ V⁻¹). The electronic noise spectral density in transparent wafers is much lower as illustrated in Figure S8 (Supporting Information). The trap density (n_{traps}) of the MAPbBr₃ wafers was

investigated by the space-charge-limited current (SCLC) method. When the MAPbBr₃ wafer changes from opaque to transparent, the trap density decreases from 6.7×10^{10} to 1.0×10^{10} cm⁻³, approaching the value of 7.6×10^9 for single crystals (Figure 2g). The diminished defect density and improved carrier mobility of the transparent wafer could be attributed to the substantially reduced porosity and grain boundary density. Figure 2h displays their *I*–*V* curve. It is inferred that ohmic contacts formed for transparent MAPbBr₃ wafers and single crystals. All the results confirm the optoelectronic properties in our transparent MAPbBr₃ wafer are comparable to those of single crystals.

To further identify the universality of our manufacturing technique, we apply it to various perovskites. Figure S9 (Supporting Information) shows photographs of as-fabricated perovskite transparent wafers with different A-site cations, B-site cation-doped perovskites, and different halogen perovskites. All the perovskite wafers display good optical transparency, indicating the advanced nature and universality of fine ceramic manufacturing engineering in the preparation of transparent perovskite wafers,

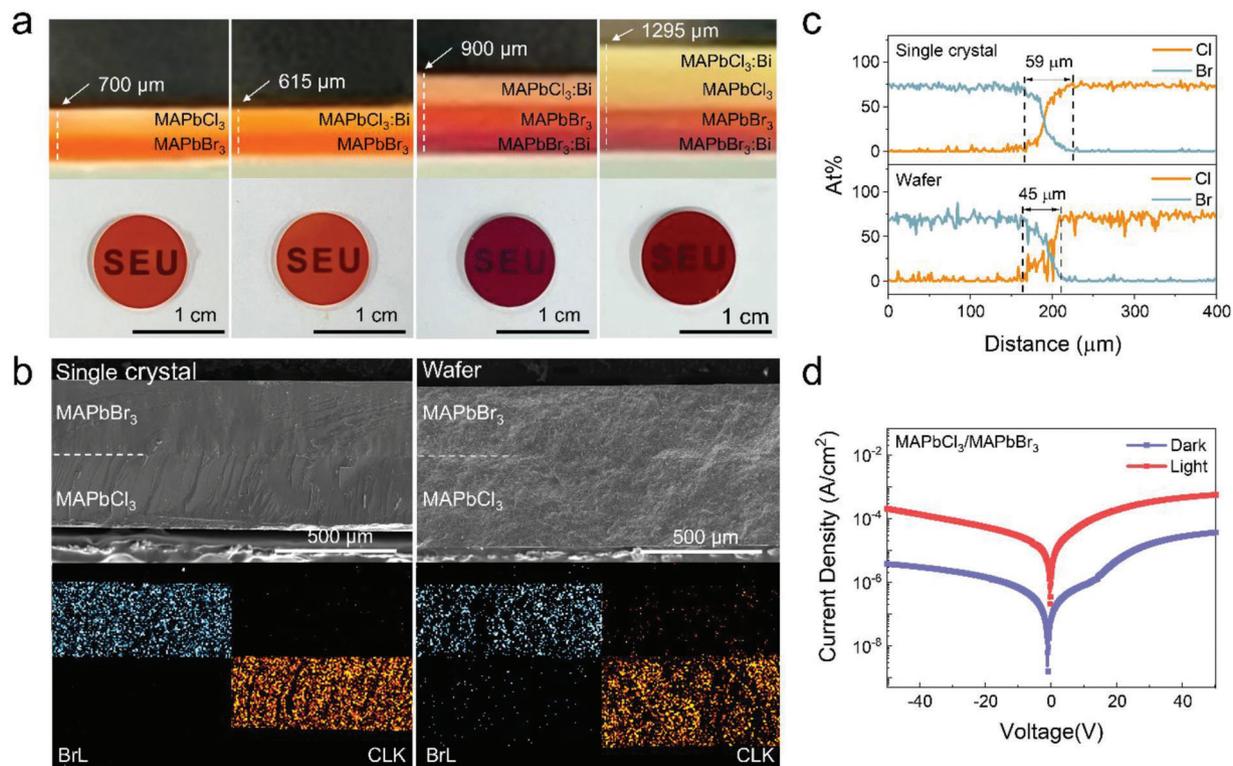


Figure 3. Characterization of the perovskite heterojunction wafers. a) Top optical and cross-sectional optical images of multilayer perovskite wafers. b) SEM images and EDS elemental mappings of MAPbCl₃/MAPbBr₃ single crystal heterojunction and heterojunction wafer. c) EDS line scan profiles of the MAPbCl₃/MAPbBr₃ heterojunction. d) *I*-*V* curves of Au/MAPbCl₃-MAPbBr₃/Au heterojunction wafer under dark and 525 nm illumination conditions.

which lay a foundation for the industrial fabrication of perovskite-based devices.^[23]

Even more, the transparent perovskite wafer manufacturing technique can also be extended to prepare double-layer and even multilayer perovskite heterojunction wafers. **Figure 3a** presents the photographs of our as-fabricated perovskite heterojunction wafers, including MAPbCl₃/MAPbBr₃, MAPbCl₃:Bi/MAPbBr₃, MAPbCl₃:Bi/MAPbBr₃:Bi and MAPbCl₃:Bi/MAPbCl₃/MAPbBr₃:Bi (The transparent perovskite wafer manufacturing technique is universal for different types of perovskite heterojunctions, taking Bi ion doping as an example). In **Figure 3b**, the cross-section SEM gives the detailed microstructures of transparent MAPbCl₃/MAPbBr₃ heterojunction wafers in comparison with the epitaxial single crystal heterojunction. The wafer layers exhibit uniform thickness distribution, and the heterojunction shows excellent uniformity along with compact interlayer connections and clear boundaries. The energy-dispersive spectroscopy (EDS) mapping analysis reveals a clear border between Chlorine (Cl) and bromine (Br) across different layers, indicating the distinct compositional characteristics of each layer, further corroborating the heterogeneous nature of the material. However, a halide ion diffusion length of 45 μm was also observed in the quantitative EDS line-scan analysis (**Figure 3c**), which can be attributed to halide ion diffusion. Even so, the diffusion region is also significantly smaller than that in the single-crystal heterojunction (59 μm). This result further confirms that our perovskite wafers can effectively suppress ion diffusion. The photoluminescence

(PL) properties of the MAPbCl₃/MAPbBr₃ heterojunction wafers were studied. As shown in **Figure S10** (Supporting Information), the heterojunction wafer revealed two well-resolved emission peaks corresponding to the constituent MAPbCl₃ and MAPbBr₃ phases. The absence of intermediate emission features between the characteristic MAPbCl₃ and MAPbBr₃ PL peaks indicates the formation of a built-in electric field at the heterointerface, which facilitates efficient transport of photogenerated charge carriers within the heterojunction architecture. The energy-band alignment at the heterojunction plays a decisive role in the separation and extraction of photogenerated charge carriers. Based on the ultraviolet photoelectron spectroscopy (UPS) (**Figure S11a**, Supporting Information) and absorption spectra (**Figure S11b,c**, Supporting Information), the valence band maximum (E_{VBM}), Fermi level (E_{F}), and conduction band minimum (E_{CBM}) of the wafers were calculated, as shown in **Figure S11d**. The E_{F} of the MAPbCl₃ wafer is close to the E_{CBM} , indicating its nature as an *n*-type semiconductor, whereas the E_{F} of the MAPbBr₃ wafer is near the E_{VBM} , suggesting its classification as a *p*-type semiconductor. With the band alignment of MAPbCl₃ and MAPbBr₃ in thermodynamic equilibrium, an internal electric field from *n*-type MAPbCl₃ to *p*-type MAPbBr₃ is formed, as shown in **Figure S12**, and the field significantly affected the efficient carrier separation. To further investigate the electron transport properties, the current density–voltage (*I*-*V*) curves of heterojunction wafer under dark and 525 nm illumination were measured (**Figure 3d**), which exhibited a typical diode effect, namely, conduction at positive bias and cutoff at negative bias,

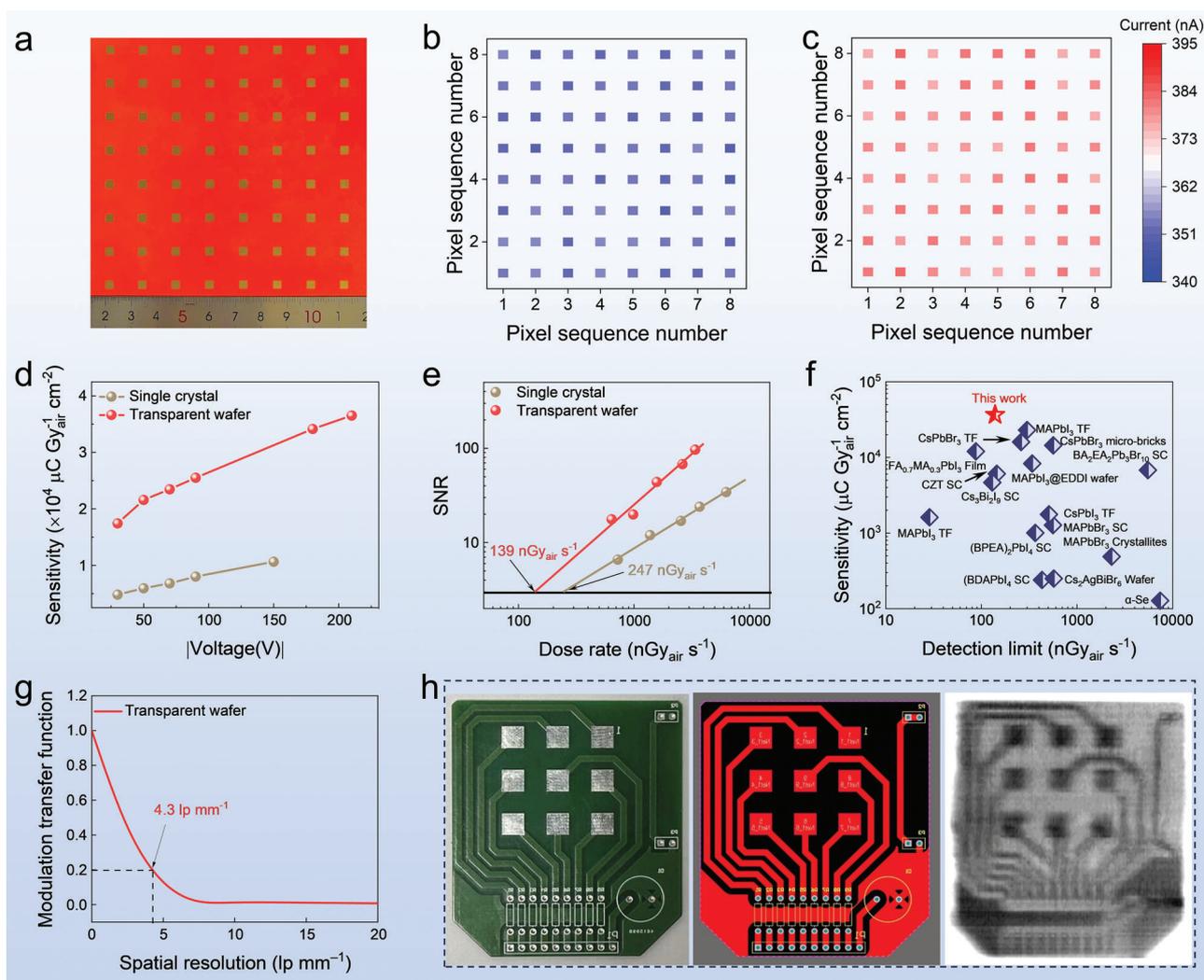


Figure 4. X-ray imaging based on the transparent perovskite heterojunction wafer. a) An optical image of an 8×8 pixels array fabricated from the transparent wafer, with a pixel size of 3 mm and a pitch size of 10 mm over a device area of $10 \times 10 \text{ cm}^2$. The current uniformity of pixels (area of each pixel: 3 mm \times 3 mm) as the regions of b) dark current and c) photocurrent. d) X-ray sensitivity. e) SNR of the pixels. f) The comparison of sensitivity and detection limit with previously reported perovskite X-ray detectors. g) The spatial resolution of the transparent wafer X-ray sensing array contains 256×256 pixels. h) Optical image, schematic diagram, and X-ray image of a printed circuit board (PCB) with an array of 256×256 pixels.

with distinct current values under different biases. In particular, under a negative bias, owing to the matching directions of the applied electric field and the built-in electric field in the space charge region, the energy-band barrier of the heterojunction was significantly increased, which effectively prevents electrons and holes from migrating through the space-charge region, and thus contributed to a decrease in the dark current.

To assess the potential for practical application, we employ our transparent perovskite wafer for constructing functional devices. As X-ray sensing sets stringent requirements on active material, we chose the X-ray sensing array on a transparent MAPbCl₃/MAPbBr heterojunction wafer. Figure S13 (Supporting Information) shows its attenuation efficiency for 40 keV X-ray photons. The MAPbBr₃ wafer exhibits strong X-ray attenuation due to its high absorption coefficient. As shown in Figure 4a, we first established an 8×8 -pixel array was established on a heterojunction wafer with an area of $10 \times 10 \text{ cm}^2$. To evaluate the uni-

formity of pixels, we measured the dark current and photocurrent of each pixel (under an X-ray dose rate of $20 \mu\text{Gy}_{\text{air}} \text{ s}^{-1}$) at a bias voltage of -30 V . As shown in Figure 4b,c, the standard deviation to average signal ratio was calculated to be 5.3% for dark current and 8.4% for photocurrent, respectively. Both are better than standard regulations.^[14b] Figure 4d presents the sensitivity under various electric fields, which was derived from the slope of the photocurrent densities versus X-ray dose rates (Figures S14–S16, Supporting Information). As the working voltage increases (from -30 to -210 V), the sensitivity of our transparent wafer device rises from 1.74×10^4 to $3.65 \times 10^4 \mu\text{C Gy}_{\text{air}}^{-1} \text{ cm}^{-2}$, which is ≈ 4 times higher than that of the single crystal device. Figure 4e shows the detection limit under a signal-to-noise ratio (SNR) of 3. Under the applied voltage of -30 V , the detection limit was estimated to be $139 \text{ nGy}_{\text{air}} \text{ s}^{-1}$, which is better than that of the single crystal device ($247 \text{ nGy}_{\text{air}} \text{ s}^{-1}$). Notably, The obtained sensitivity is three orders of magnitude

higher than that of α -Se ($20 \mu\text{C Gy}_{\text{air}}^{-1} \text{cm}^{-2}$),^[24] five times higher than that of the CZT detector ($6069 \mu\text{C Gy}_{\text{air}}^{-1} \text{cm}^{-2}$)^[8] and much higher than that of MAPbI₃ films ($23\,000 \mu\text{C Gy}_{\text{air}}^{-1} \text{cm}^{-2}$),^[25] CsPbBr₃ film ($1450 \mu\text{C Gy}_{\text{air}}^{-1} \text{cm}^{-2}$)^[26] and Cs₂AgBiBr₆ Wafer ($250 \mu\text{C Gy}_{\text{air}}^{-1} \text{cm}^{-2}$)^[14b] (Figure 4f; Table S5, Supporting Information). The detection limit is ≈ 38 times lower than the current medical diagnostic threshold ($5.5 \mu\text{Gy}_{\text{air}} \text{s}^{-1}$) and comparable to the reported representative single crystal detectors (Figure 4f; Table S5, Supporting Information).^[27] Figure S17 (Supporting Information) shows faster rise and fall times (24.7 and 25.4 μs) in the pixel, enabling high-speed X-ray imaging. The response stability of the MAPbCl₃/MAPbBr₃ heterojunction wafer detector was measured under X-rays with a dose rate of $95 \mu\text{Gy}_{\text{air}} \text{s}^{-1}$ and at a bias of -30 V for 3700 s, as shown in Figure S18a (Supporting Information). It can be seen that both the photocurrent and dark current of the detector remain stable without obvious degradation throughout the measurement period. The long-term stability of the heterojunction wafer in an air atmosphere is the key characterization for their practical applications. The response current of the wafer detector was calculated and remained at 97% of the initial value after 15 days (Figure S18b, Supporting Information). The excellent stability of the detector is attributed to the high crystalline quality and suppressed ion migration of the MAPbCl₃/MAPbBr₃ heterojunction wafer.

According to the established regulations, the artifacts and the ratio of the standard deviation to the mean signal in X-ray imaging should be less than 10%.^[14b] Thus, we further fabricated a sensing array containing 256×256 pixels on a $10 \times 10 \text{ cm}^2$ transparent wafer with a pixel pitch of $200 \mu\text{m}$. By using the slanted-edge method,^[22] the spatial resolution of the sensing array was calculated to be 4.3 lp mm^{-1} when the modulation transfer function of the detector was 0.2 (Figure 4g; Figure S19, Supporting Information). The object used for array imaging, as illustrated in Figure 4h and Figure S20 (Supporting Information), consists of a circuit board featuring diverse materials, including an epoxy resin substrate and copper circuit traces. During imaging, the X-ray dose rate was set to that of regular medical diagnostics ($5.5 \mu\text{Gy s}^{-1}$). The images contain a clear and comprehensive morphology and internal structure of the circuit board. Additionally, discrimination among materials of different compositions has been achieved. It confirms that our sensing array based on a transparent wafer is suitable for high-quality X-ray imaging.

3. Conclusion

In summary, this work presents a breakthrough in scalable manufacturing of wafer-scale lead halide perovskite wafers, achieving optoelectronic properties approaching those of single crystals with the ceramic manufacturing technique. By precisely controlling temperature, pressure, and duration during fabrication, we successfully produced 6-inch transparent MAPbBr₃ wafers with exceptional optical transparency ($\approx 68\%$ in 700–1100 nm), high carrier mobility (hole: $99 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$; electron: $180 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$), and ultralow trap density ($1.0 \times 10^{10} \text{ cm}^{-3}$), rivaling the performance of MAPbBr₃ single crystals. The technique demonstrates universality across diverse perovskite compositions (A/B/X-site variations) and multilayer heterojunctions. Leveraging these advances, we fabricated a $10 \times 10 \text{ cm}^2$ X-ray sensing array with 256×256 pixels, achieving ultrahigh sensitivity ($36532 \mu\text{C Gy}_{\text{air}}^{-1}$

cm^{-2}) and ultralow detection limit ($139 \text{ nGy}_{\text{air}} \text{ s}^{-1}$), surpassing both single-crystal detectors and commercial counterparts. We believe this work bridges the critical gap between lab-scale perovskite research and industrial manufacturing, offering a universal platform for wafer-scale integration of perovskite devices with silicon-compatible processes.

4. Experimental Section

Preparation of the MAPbBr₃ Perovskite Wafers: Perovskite microcrystals were synthesized using the inverse temperature crystallization method, followed by drying and grinding into finer powder particles.^[28] The specific procedure entailed the transfer of a 15 mL solution of N,N-dimethylformamide (DMF) into a beaker using a pipette. Subsequently, 3.54 g of methylammonium bromide (MABr) powder was added to the beaker and stirred on a magnetic stirrer until the salt was completely dissolved. Then, 11.12 g of lead bromide (PbBr₂) powder was added and stirred until the solution became clear and transparent. The solution was then filtered using a $22 \mu\text{m}$ organic filter to prepare a precursor solution with an MABr concentration of 1 mol L^{-1} . The precursor solution was placed on a stirring heater and stirred for 3 h at $80 \text{ }^\circ\text{C}$, leading to the precipitation of MAPbBr₃ microcrystals as the solubility decreased. The microcrystals were subsequently dried in a vacuum dryer at $60 \text{ }^\circ\text{C}$ for 12 h, followed by grinding into smaller particles ($\approx 50 \mu\text{m}$) using an agate mortar and pestle to obtain the MAPbBr₃ perovskite powder necessary for hot pressing. The powder was then loaded into a hot-pressing mold, and the corresponding temperature, pressure, and time settings were adjusted before applying pressure once the set temperature was reached. Following hot pressing, the equipment was allowed to cool down to $50 \text{ }^\circ\text{C}$, and the hot-pressed wafers were subsequently removed from the mold.

Preparation of Perovskite Heterostructures Wafers: The Perovskite heterojunction wafers were fabricated via a layer-by-layer hot-press assembly. The MAPbBr₃ microcrystalline precursor powders were sequentially loaded into a custom hot-press mold. An initial pressure of $\approx 10 \text{ MPa}$ was applied after the deposition of each precursor layer to ensure uniform thickness, complete coverage, and homogeneous powder distribution. The subsequent layer (the MAPbCl₃ microcrystalline precursor powders) was added only after the preceding layer (MAPbBr₃) stabilized. After complete assembly, set the appropriate temperature and time, and immediately apply pressure to the hot-press mold once the set temperature was reached. Following hot pressing, the equipment was allowed to cool down to $50 \text{ }^\circ\text{C}$, and the hot-pressed heterojunction wafers were subsequently removed from the mold. For multilayer perovskite heterostructures, an initial pressure of $\approx 10 \text{ MPa}$ was applied after the deposition of each precursor layer to ensure uniform thickness, complete coverage, and homogeneous powder distribution. Subsequent layers were added only after the preceding layer stabilized. After complete assembly, the aforementioned hot pressing process was repeated, producing monolithic perovskite heterocrystalline wafers. The wafer dimensions and thickness were controlled by adjusting the mold size and the mass of precursor powder.

Fabrication of the Detectors: A vertical Au-MAPbBr₃-Au photodetector structure was fabricated through a process involving the application of a rigid masking film on both the top and bottom surfaces of the MAPbBr₃ wafer, followed by the deposition of square Au electrodes (50 nm thick, with an area of $2.5 \text{ mm} \times 3 \text{ mm}$) within a vacuum chamber.

In the case of the sensing array, it was fabricated on the transparent perovskite heterojunction wafer with a circular area measuring 15 cm in diameter. The top and bottom surfaces of the polycrystalline wafer were then affixed to PCB boards featuring horizontally and vertically arranged electrodes, utilizing a specially formulated crystal adhesive developed within the research group. The overlapping regions of the electrodes served as pixel points (size of $200 \mu\text{m} \times 200 \mu\text{m}$). The preparation of the crystal adhesive entailed a precise procedure, involving the addition of ethyl cellulose powder (1 g each, added thrice) into 9 mL of turpentine alcohol. The resulting mixture was then stirred and heated on a magnetic stirrer

heating plate at 90 °C for 2 h. Following this, 10 mL of MAPbBr₃ precursor solution was added to a beaker and vigorously stirred with a glass rod to promote fusion, after which the solution was further stirred and heated for 1 h to yield an orange viscous crystal adhesive.

Characterizations: Transmission spectra of all samples were acquired using a Cary 5000 UV–vis–NIR spectrophotometer (Agilent Technologies). The steady-state PL spectra were collected with an FS5 spectrofluorometer (Edinburgh Instruments, UK). The TRPL measurements were performed using a picosecond pulsed diode laser for excitation and a time-correlated single-photon counting detector for signal collection. The X-ray diffraction patterns were tested using a Rigaku SmartLab SE X-ray diffractometer and Cu K α radiation (1.5406 Å). SEM images and EDS were investigated using Field-Emission scanning electron microscopy (Quanta 200 FEI). The noise current was extracted from the dark current recorded by an Agilent oscilloscope (7.5G) with a current amplifier. The electrical characteristics, represented by I - V and I - T curves, were measured using a Keithley 2400 parameter analyzer. To determine the rise time and fall time, the frequency of the light source (525 nm LED) was controlled utilizing an Agilent 33500B function generator. X-ray was generated by an Amptek Mini-X2 and the dose rate was calibrated using a 10 \times 6-180 dosimeter (Radcal, USA). For X-ray imaging purposes, the detector connected to the PCB board was interfaced with the array imaging system developed by the research team via a ribbon cable, facilitating amplification and readout of pixel charges.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

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Conflict of Interest

The authors declare no conflict of interest.

Data Availability Statement

The data that support the findings of this study are available in the supplementary material of this article.

Keywords

ceramic, heterojunction, perovskites wafer, six-inch, X-ray array

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